



CLASSICAL SOLUTIONS OF THE GENERALIZED CAMASSA-HOLM EQUATION



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ABSTRACT

In this poster, well-posedness in $C^1(\mathbb{R})$ (a.k.a. classical solutions) for a generalized Camassa-Holm equation (g - kb CH) having $(k+1)$ -degree nonlinearities is explored. This result holds for the Camassa-Holm, the Degasperi-Procesi and the Novikov equations, which improves upon earlier results in Sobolev and Besov spaces.

MAIN INGREDIENTS

For $k \in \mathbb{Z}^+$ and $b \in \mathbb{R}$, we consider the Cauchy problem for the following generalized Camassa-Holm (g - kb CH) equation

$$\begin{cases} (1 - \partial_x^2) \partial_t u = u^k \partial_x^3 u + bu^{k-1} \partial_x u \partial_x^2 u - (b+1)u^k \partial_x u, \\ u(x, 0) = u_0(x), \quad x \in \mathbb{R} \text{ and } t \in \mathbb{R}, \end{cases}$$

which takes the following non-local form

$$\partial_t u + u^k \partial_x u + F(u) = 0,$$

$$\begin{aligned} F(u) = & \partial_x (1 - \partial_x^2)^{-1} \left[\frac{b}{k+1} u^{k+1} + \frac{3k-b}{2} u^{k-1} (\partial_x u)^2 \right] \\ & + (1 - \partial_x^2)^{-1} \left[\frac{(k-1)(b-k)}{2} u^{k-2} (\partial_x u)^3 \right]. \end{aligned}$$

Theorem 1 (Picard-Lindelöf) *Let X be a Banach space. Suppose that $f : X \rightarrow X$ is locally Lipschitz on a closed ball $\bar{B}_R(u_0) \subset X$ where $R > 0$ and $u_0 \in X$. Let*

$$M = \sup_{u \in \bar{B}_R(u_0)} \|f(u)\| < \infty.$$

Then the initial value problem

$$\begin{cases} \dot{u} = f(u) \\ u(0) = u_0 \end{cases}$$

has a continuously differentiable local solution $u(t)$. This solution is defined in the time interval $t \in (-\delta, \delta)$ where $\delta = R/M$.

REFERENCES

- J. Holmes, R. Thompson, *Classical Solutions for the Generalized Camassa-Holm Equation*, Submitted to *Advances in Differential Equations*, (2015).
- A. Himonas, R. Thompson. *Persistence properties and unique continuation for a generalized Camassa-Holm equation*, *Journal of Mathematical Physics*, 55 091503 (2014).

OBJECTIVE

We will show that this family of shallow water wave equations is well-posed in the space of bounded and continuously differentiable functions on the real line, denoted C^1 , and equipped with the norm

$$\|f\|_{C^1} = \sup_{x \in \mathbb{R}} |f(x)| + \sup_{x \in \mathbb{R}} \left| \frac{d}{dx} f(x) \right|.$$

METHODOLOGY

We begin by showing how one formally constructs an equivalent ODE system to the g - kb CH equation. Assuming a solution, u , exists and is a C^1 solution of the g - kb CH initial value problem, we have our trajectories satisfy the ODE

$$\begin{cases} \eta_t(x, t) = u^k(\eta(x, t), t) \\ \eta(x, 0) = x. \end{cases}$$

Moreover, the above ODE has a unique solution $\eta(x, t)$ which is also continuously differentiable, therefore, we may define

$$\begin{aligned} w(x, t) &= u(\eta(x, t), t), \quad v(x, t) = u_x(\eta(x, t), t), \\ q(x, t) &= \eta_x(x, t), \end{aligned} \quad (3)$$

and we see that we may easily obtain $u(x, t)$ from the composition

$$u = w \circ \eta^{-1}.$$

We will first find a system of equations satisfied by w , v and q , and then show that this system of equations is indeed an ODE system, and therefore the solutions are uniquely defined. Using w , we will then construct η and u similarly to the above formal definitions.

FUTURE RESEARCH

- Classical solutions to the CH2 system and other shallow water wave equations.
- Asymptotic profiles and propagation speed of solutions to other shallow water wave systems.
- Local and global solutions to Camassa-Holm type equations in Besov spaces.

THE SEMI-LINEAR SYSTEM

$$\begin{cases} \frac{d}{dt} w = -P_1(w, v, q) - R_1(w, v, q), \\ \frac{d}{dt} v = \frac{k-b}{2} w^{k-1} v^2 + \frac{b}{k+1} w^{k+1} - P_2(w, v, q) - R_2(w, v, q), \\ \frac{d}{dt} q = kw^{k-1} vq, \end{cases} \quad (1)$$

with initial data

$$\begin{cases} w(x, 0) = u_0(x) \\ v(x, 0) = \frac{d}{dx} u_0(x) \\ q(x, 0) = 1, \end{cases} \quad (2)$$

where

$$\begin{aligned} P_1(w, v, q) & \doteq \frac{1}{2} \int_x^\infty e^{-|\int_x^z q(y,t) dy|} \left(\frac{b}{k+1} w^{k+1} q + \frac{3k-b}{2} w^{k-1} v^2 q \right) (z, t) dz \\ & - \frac{1}{2} \int_{-\infty}^x e^{-|\int_x^z q(y,t) dy|} \left(\frac{b}{k+1} w^{k+1} q + \frac{3k-b}{2} w^{k-1} v^2 q \right) (z, t) dz. \end{aligned}$$

$$P_2(w, v, q) = \frac{1}{2} \int e^{-|\int_x^z q(y,t) dy|} \left[\frac{b}{k+1} w^{k+1} + \frac{3k-b}{2} w^{k-1} v^2 \right] q dz$$

$$R_1(w, v, q) \doteq \frac{(k-1)(b-k)}{4} \int e^{-|\int_x^z q(y,t) dy|} w^{k-2}(z, t) v^3(z, t) q(z, t) dz.$$

$$\begin{aligned} R_2(w, v, q) & \doteq \frac{(k-1)(b-k)}{2} \frac{1}{2} \int_x^\infty e^{-|\int_x^z q(y,t) dy|} w^{k-2}(z, t) v^3(z, t) q(z, t) dz \\ & - \frac{(k-1)(b-k)}{2} \frac{1}{2} \int_{-\infty}^x e^{-|\int_x^z q(y,t) dy|} w^{k-2}(z, t) v^3(z, t) q(z, t) dz. \end{aligned}$$

SUMMARY OF PROOF FOR WELL-POSEDNESS

- We show that the forcing terms from the o.d.e. (1) are locally Lipschitz in the space $C^1 \times C \times C$.
- This gives us a unique solution (w, v, q) within the time interval $t \in [-T, T]$ where $T = \min \left\{ \frac{1}{2\|u_0\|_{C^1}^k}, \frac{1}{2L} \right\}$, where L is our Lipschitz constant.
- We construct our solution $u(x, t)$ from our particle trajectories $\eta : \mathbb{R} \rightarrow \mathbb{R}$ and show uniqueness and continuous dependence on the initial data.

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